

### 7.2.3 Response of a Narrowband Filter to Noise

If the representation of noise as a superposition of spectral components is a reasonable one, we should expect that when noise is passed through a narrowband filter the output of the filter should look rather like a sinusoid. We find that such is indeed the case, for we find that the output of a narrowband filter with noise input has the appearance shown in Fig. 7.4. The output waveform looks like a sinusoid except that, as expected, the amplitude varies randomly. The spectral range of the envelope of the filter output encompasses the spectral range from  $-B/2$  to  $B/2$ , where  $B$  is the filter bandwidth. The average frequency of the waveform is the center frequency  $f_c$  of the filter. If  $B \ll f_c$ , the envelope changes very "slowly" and makes an appreciable change only over many cycles. Thus, while the spacings of the zero crossings of the waveform are not precisely constant, the change from cycle to cycle is small and when averaged over many cycles is quite constant at the value  $1/2f_c$ . Finally, we may note that as  $B$  becomes progressively smaller, so also does the average amplitude, and the waveform becomes more and more sinusoidal.

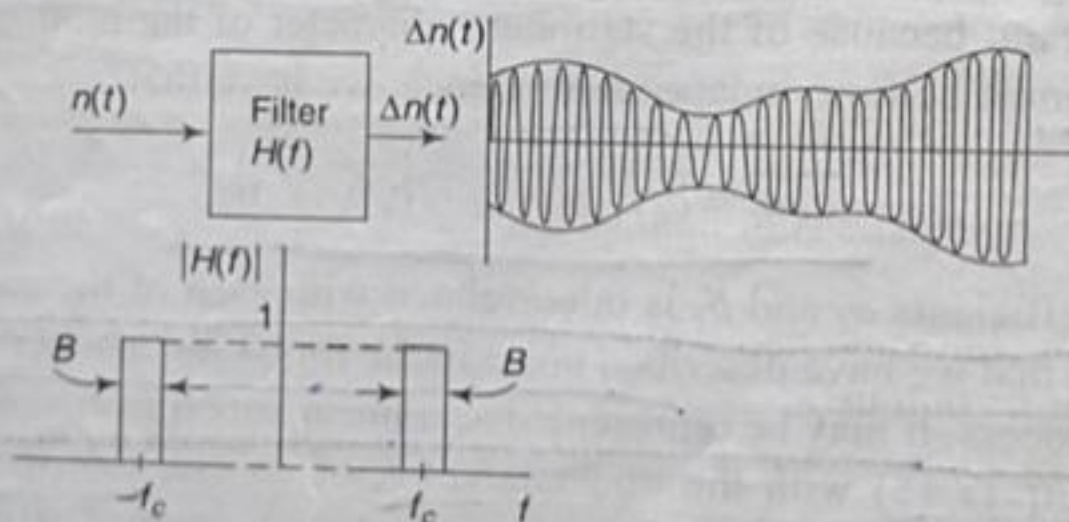


Fig. 7.4 Response of a narrowband filter to noise.

### 7.2.4 Effect of a Filter on PSD of Noise

Let a spectral component of noise  $n_k(t)$  given by Eq. (7.15) be the input to a filter whose transfer function at the frequency  $k \Delta f$  is

$$H(k \Delta f) = |H(k \Delta f)| e^{j\varphi_k} = |H(k \Delta f)| / \varphi_k \quad (7.29)$$

The corresponding output spectral component of noise will be  $n_{k_o}(t)$

$$n_{k_o}(t) = |H(k \Delta f)| a_k \cos(2\pi k \Delta f t + \varphi_k) + |H(k \Delta f)| b_k \sin(2\pi k \Delta f t + \varphi_k) \quad (7.30)$$

The power  $P_k$  associated with  $n_k(t)$  is, from Eq. (7.20),

$$P_k = \frac{\overline{a_k^2} + \overline{b_k^2}}{2} \quad (7.31)$$

Since  $|H(k \Delta f)|$  is a deterministic function,  $[|H(k \Delta f)|a_k]^2 = |H(k \Delta f)|^2 \overline{a_k^2}$ , and  $[|H(k \Delta f)|b_k]^2 = |H(k \Delta f)|^2 \overline{b_k^2}$ . Hence, comparing Eq. (7.30) with Eq. (7.15), we find that the power  $P_k$  associated with  $n_k(t)$  is

$$P_k = |H(k \Delta f)|^2 \frac{\overline{a_k^2} + \overline{b_k^2}}{2} \quad (7.32)$$

Finally, then, from Eqs (7.31) and (7.32), using also Eq. (7.20), we have that the power spectral densities at input and output,  $G_{n_i}(k \Delta f)$  and  $G_{n_o}(k \Delta f)$ , are related by

$$G_{n_o}(k \Delta f) = |H(k \Delta f)|^2 G_{n_i}(k \Delta f) \quad (7.33)$$

In the limit as  $\Delta f \rightarrow 0$  and  $k \Delta f$  is replaced by a continuous variable  $f$ , Eq. (7.33) becomes

$$G_{n_o}(f) = |H(f)|^2 G_{n_i}(f) \quad (7.34)$$

Note the similarity between this result and Eq. (1.81), which applies to a deterministic waveform.

We got this relation by convolution for a transmission of a random process through a linear system in Sec. 6.5.6.